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Remote Measurement Using a Femtosecond-Optical-Frequency-Comb-Based Interferometer

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Abstract: A novel interferometric technique using a femtosecond optical frequency comb (FOFC) is developed to measure the depth difference between two gauge blocks. The present technique is expected to be useful for remote measurement.

Keywords: Mode-locked lasers, Temporal Coherence, Interferometry, Remote measurement

1. INTRODUCTION

The increased frequency stability and the very-broad frequency band of femtosecond optical frequency comb (FOFC) has led to the application of this device in several precision metrology areas such as high-precision optical metrology, spectroscopy, and measurements. We have investigated the temporal coherence function (TCF) of a pulse train from the FOFC [1]. The results show that high temporal coherence peaks exist during the period equal to the repetition intervals in the traveling direction of the FOFC. In the present study, we demonstrated an unbalanced optical-path Michelson interferometer for remote measurement by observation of high temporal coherence between different pairs of pulse trains from the FOFC. As a special demonstration, measurement of the depth difference of two gauge blocks with 3 m unbalanced optical path was demonstrated. For simplicity of explanation, we have neglected the dispersion and absorption of the optical elements over the FOFC's illumination bandwidth.

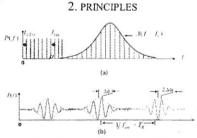


Fig. 1. Fourier-transform relationship between the power spectrum and the interference fringes. (a) Power spectrum of the FOFC. (b) Interference fringes; dotted line, the temporal coherence function of the FOFC; solid line: interference fringes.

In what follows, for convenience of explanation, let us first consider the temporal coherence function of the FOFC. The power spectrum of an FOFC light source can be expressed as:

$$P(f) \propto A(f - f_c) \times \sum_{m=-\infty}^{\infty} \delta(f - mf_{rep} - f_{CEO}),$$
 (1)

where $A(f-f_{\rm c})$ is the envelope function of the FOFC power spectrum, $f_{\rm c}$ is the center carrier frequency of the FOFC. When the electric field packet repeats at the pulse repetition period $T_{\rm R}$, the "carrier" phase slips by $\Delta\varphi_{\rm u}$ to the carrier-envelope phase because of the difference between the group and phase velocities. In the frequency domain, a mode-locked FOFC generates equidistant frequency comb lines with the pulse repetition frequency $f_{\rm rep} \propto 1/T_{\rm R}$, and due to phase slip $\Delta\varphi_{\rm ce}$, the whole equidistant-frequency comb is shifted by $f_{\rm CEO}$. Based on the Wiener-Khintchine theorem, the interferometric signal of the autocorrelation function is given by

$$\gamma(\tau) \propto F^{-1} \left[A(f - f_{\rm c}) \right] \otimes \sum_{m=-\infty}^{+\infty} \delta(\tau - mT_{\rm R}).$$
 (2)

From Eq. (2), the temporal coherence function periodically displays a high temporal coherence peak where the pulse trains signal of the FOFC displays a high-intensity peak with the pulse repetition period T_{ν} .

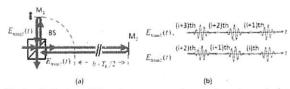


Fig.2. Relative delay between pulse trains formed by an unbalanced optical-path Michelson interferometer. (a) Simplified optical layout for interferometry. (b) Relative positions between two pulse trains (integer h=1).

Let us consider the interference fringes formed by an unbalanced optical-path Michelson interferometer, as shown in Fig. 2(a). First, the incoming pulse train is split into two identical parts. The mode-lock technique results in interference fringes reappearing at delays equal to $hT_{\rm g}$. When the pulse trains $E_{\rm main1}(t)$ and the relatively delayed pulse trains $E_{\rm main2}(t)$ finally overlap at the BS (see Fig. 2(b)), one can expect that interference fringes can be observed, as shown in Fig. 1(b). After performing the time integration we obtain

$$I(t) \propto |\gamma(\tau)| \cos[\operatorname{mod}(h \times \Delta \varphi_{cc}, 2\pi)].$$
 (3)

3. EXPERIMENT

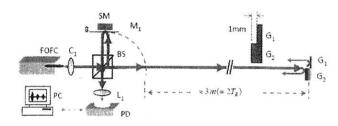


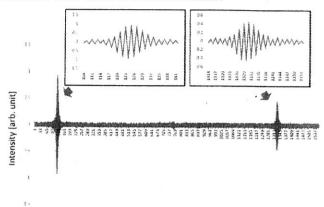
Fig. 3. Basic scheme for experiment. Abbreviations defined in text.

The basic scheme of the experiments is shown in Fig. 3. The experiment is carried out with a system consisting of an FOFC (FC1500, Menlo Systems, $f_{rep} = 100 \,\text{MHz}$), an unbalanced optical-path Michelson interferometer, and system controls. The pulse trains from the FOFC is expanded and collimated by a collimator C1 and introduced into the unbalanced optical-path Michelson interferometer. The unbalanced optical-path Michelson interferometer composed of a beam splitter BS, a reference mirror M1, and a combination object with two gauge blocks G1 and G2 with depth difference 1 mm. During the measurement, by moving the reference mirror M₁ of the interferometer by means of a computer-controlled and calibrated ultrasonic stepping motor SM, we could observe the interference fringes. After travelling different path lengths, these two pairs of pulse trains overlap at the BS. Lens L1 images the interference fringes onto a photo detector PD.

Figure 4 illustrates the acquire interference fringes. The 1st and 2rd interference fringe signals exhibit a high contrast between the two pairs of pulse trains by the relative unbalanced optical displacement 2cT_R (about 3 m, c is the light velocity in air.).

To obtain the relative average depth difference between two gauge blocks we analyzed the fringes pattern in Fig.4. Because the sampling theorem did not satisfy as noted in Fig.4, the Fourier transform method [1] was not able to be applied. The relative depth difference of two gauge blocks was measured as 939µm, a product of length corresponding to one scanning step and all the numbers of scanning steps

between two peaks of fringes patterns. At a constant temperature, the fluctuations of displacement between two peaks of fringes patterns due to noise, such as mechanical vibration in the state, was measured as $2\mu m$. The relative average depth difference of two gauge blocks was measured as $939\pm2\mu m$. The experimental result is in agreement with the theoretical value within a standard deviation.



Sequential number of scan steps n [1 step = 775 nm]

Fig. 4. Measured interference fringes.

4. SUMMARY AND FUTURE WORK

In summary, we have presented a novel interferometry technique using an FOFC to measure the depth difference of two gauge blocks with unbalanced optical-path about 3 m. The present technique is expected to be useful for not only surface profilometry but also tomography.

Note that in these experiments, the accuracy of the measurement is limited to the restricted resolution of the digital oscilloscope. We are currently pursuing an approach in which performance can be improved using an optical component with high accuracy and an appropriate design of the system, such as high-speed measurement.

ACKNOWLEDGMENTS

Part of this research work was supported by the Global Center of Excellence (COE) Program on "Global Center of Excellence for Mechanical Systems Innovation" granted to The University of Tokyo, from the Japanese Government. We are also grateful to NEOARK Corporation for providing the femtosecond fiber laser. Dong Wei gratefully acknowledges the scholarship given from Takayama International Education Foundation, Heiwa Nakajima Foundation, Ministry of Education, Culture, Sports, Science, and Technology of Japan, respectively.

[1] Dong Wei, Satoru Takahashi, Kiyoshi Takamasu, and Hirokazu Matsumoto, "Analysis of the temporal coherence function of a femtosecond optical frequency comb," Opt. Express 17, 7011-7018 (2009)